

Aharonov-Bohm oscillations in a mesoscopic ring with two entangled magnetic impurities

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cond-mat/0612057 (2006)



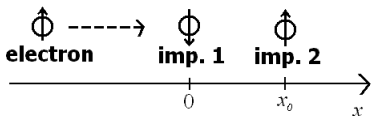
Outline

- 1 1D Wire with two entangled magnetic impurities
 - Electron coherence via entanglement of the impurities
- 2 Mesoscopic ring with two entangled magnetic impurities

1D Wire with two spin-1/2 impurities

F. Ciccarello *et al.*, New. J. Phys. **8**, 214 (2006)

F. Ciccarello *et al.*, sottomesso (2006), ArXiv:quant-ph/0611025



$$H = \frac{p^2}{2m^*} + J\sigma \cdot \mathbf{S}_1 \delta(x) + J\sigma \cdot \mathbf{S}_2 \delta(x - x_0)$$

- Multiple scattering between 1 and 2 with spin-flip
- Electron analogue of a Fabry-Perot with two "quantum" mirrors
- What happens when the impurities are in an entangled state?

Magnetic impurities

Sources of electron decoherence

- Static impurities \Rightarrow Well fixed phase shifts
 \Rightarrow Coherence preserved
- Magnetic impurities \Rightarrow Indeterminate phase shifts
 \Rightarrow Loss of coherence

S. Datta, *Electronic Transport in Mesoscopic Systems* (Cambridge Univ. Press, Cambridge, 1995)

Stern A, Aharonov Y and Imry Y 1990 *Phys. Rev. A* **41** 3436

Imry Y 1997 *Introduction to Mesoscopic Physics* (New York: Oxford Univ. Press)



Approach

- Derivation of exact stationary states
- Technique: adapted *quantum waveguide theory*
- The distance between the impurities is a parameter of the model



Outline

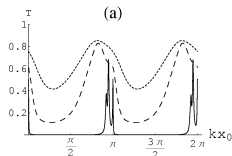
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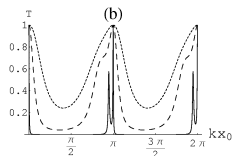
Perfect transparency

$$J/k = 1(\cdots), 2(---), 5(—)$$

- Typical case: $|\uparrow\downarrow\rangle$ (product state)



- Perfect transparency: $|\Psi^-\rangle = 2^{-1/2} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$ (particular maximally entangled state)

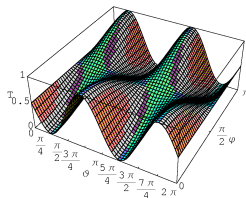


One spin-up family of states at the resonance condition

One spin-up family of impurity spin states

$$|\Psi(\vartheta, \varphi)\rangle = \cos \vartheta |\uparrow\downarrow\rangle + e^{i\varphi} \sin \vartheta |\downarrow\uparrow\rangle$$

Transmittivity vs. $\vartheta \in [0, 2\pi]$ and $\varphi \in [0, \pi]$ at $kx_0 = n\pi$



- maxima of T for $|\Psi^-\rangle = 2^{-1/2} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$
 - minima of T for $|\Psi^+\rangle = 2^{-1/2} (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)$
- F. Ciccarello *et al.*, New J. Phys. **8**, 214 (2006)



Physical explanation

- For $kx_0 = n\pi$ the electron is found at $x = 0$ and $x = x_0$ with equal probability
 $\Rightarrow \mathbf{S}_{12}^2$ additional constant of motion
- $\delta_k(\mathbf{x}) = \delta_k(x - x_0) \Rightarrow$ Effective el.-imp. coupling V :

$$V = \frac{J}{2} (\mathbf{S}^2 - \sigma^2 - \mathbf{S}_{12}^2) \delta_k(x)$$

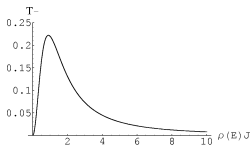
...when $|\uparrow\rangle |\Psi^-\rangle$ is prepared:

- $\mathbf{S}_{12}^2 |\uparrow\rangle |\Psi^-\rangle = 0$ $S_z |\uparrow\rangle |\Psi^-\rangle = 1/2$ (no degeneracy!)
 $\Rightarrow |\uparrow\rangle |\Psi^-\rangle \rightarrow |\uparrow\rangle |\Psi^-\rangle$: spin-flip is effectively quenched!
 $\Rightarrow V = 0$ and $H = \frac{p^2}{2m^*}$: perfect transparency!



Generation of entangled states

- for $kx_0 = n\pi$ $|\uparrow\rangle|\downarrow\downarrow\rangle \rightarrow t_\uparrow|\uparrow\rangle|\downarrow\downarrow\rangle + t_\downarrow|\downarrow\downarrow\rangle|\Psi^+\rangle$
- $T_\downarrow = |t_\downarrow|^2$: probability that the electron is transmitted in $|\downarrow\downarrow\rangle$ that is that the impurities are projected in $|\Psi^+\rangle$
- $|\Psi^+\rangle \rightarrow |\Psi^-\rangle$ through a local field acting on one of the two impurities



A. T. Costa, Jr., S. Bose, and Y. Omar, Phys. Rev. Lett. **96**, 230501 (2006)

F. Ciccarello *et al.*, New J. Phys. **8**, 214 (2006)

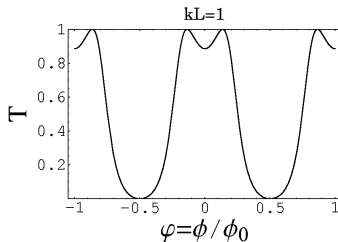
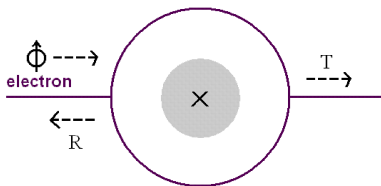
K. Yuasa, H. Nakazato, accepted for publication in J. Phys. A, ArXiv:quant-ph/0609217

G. L. Giorgi, and F. De Pasquale, Phys. Rev. B **74**, 153308 (2006)



Mesoscopic ring

Aharonov-Bohm (AB) effect (S. Washburn, and R. A. Webb, Adv. Phys. **35**, 375, 1986)

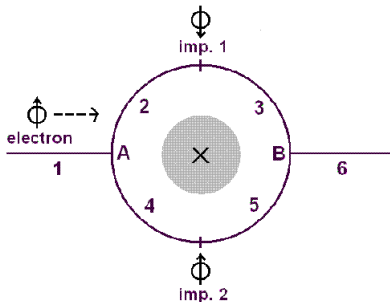


- k : incident electron wave vector
- L : circumference
- Magnetic flux: ϕ
- $\phi_0 = hc/e$: quantum flux



Mesoscopic ring with two magnetic impurities

F. Ciccarello, G. M. Palma, and M. Zarcone, submitted (2006) ArXiv: cond-mat/0612057 (2006)



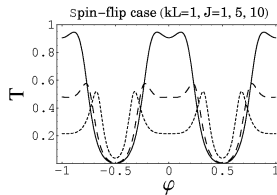
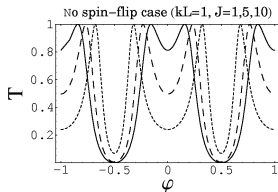
- Electron-impurity coupling: $-J \boldsymbol{\sigma} \cdot \mathbf{S}_i \delta(x_i - L/4)$
- $k_1 = k + (e\phi/\hbar cL)$, $k_2 = k - (e\phi/\hbar cL)$
- Wavefunction in the upper arm: lin. comb. of $e^{ik_1 x}$ and $e^{-ik_2 x}$
- Wavefunction in the lower arm: lin. comb. of $e^{ik_2 x}$ and $e^{-ik_1 x}$



AB effect with 1 magnetic impurity

S. K. Joshi *et al.*, Phys. Rev. B **64**, 075320 (2001)

- No spin-flip case: electron $|\uparrow\rangle$, impurity $|\uparrow\rangle$
- Spin-flip case: electron $|\uparrow\rangle$, impurity $|\downarrow\rangle$



Scattering with spin-flip \Rightarrow Loss of coherence
 \Rightarrow Reduction of the amplitude of AB oscillations

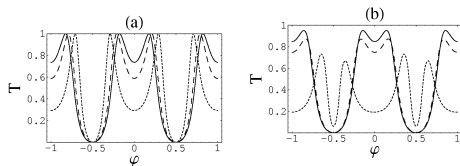


AB oscillations with 2 magnetic impurities

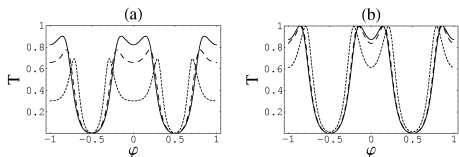
F. Ciccarello *et al.*, submitted (2006) ArXiv: cond-mat/0612057 (2006)

$kL = 1, J/k = 1$ (—), 2(---), 5(· · ·)

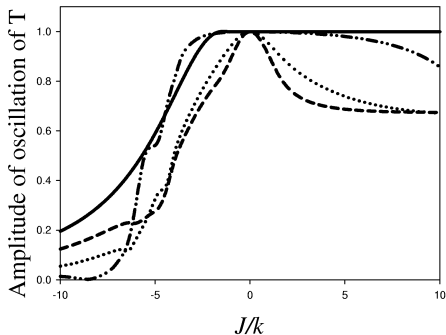
● Product states : $|\uparrow\rangle|\uparrow\uparrow\rangle$ e $|\uparrow\rangle|\uparrow\downarrow\rangle$



● Entangled states: $|\uparrow\rangle|\Psi^+\rangle$ e $|\uparrow\rangle|\Psi^-\rangle$



Amplitude of AB oscillations



$$kL = 1$$

$|\uparrow\rangle|\uparrow\rangle$ (—) J. M. Mao *et al.*, J. Appl. Phys. **73**, 1853 (1993)

$|\uparrow\rangle|\uparrow\downarrow\rangle$ (.....) Spin-flip and decoherence

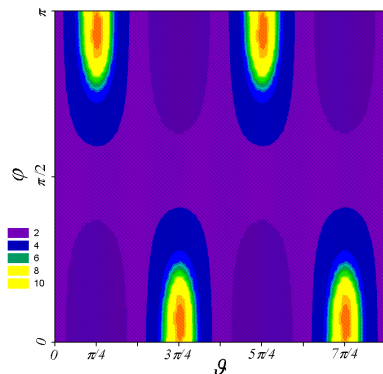
$|\uparrow\rangle|\Psi^+\rangle$ (— —) Spin-flip and decoherence

$|\uparrow\rangle|\Psi^-\rangle$ (- · · -) **Nearly coherent behaviour**

Family of one spin-up impurity spins states

$$|\Psi(\vartheta, \varphi)\rangle = \cos \vartheta |\uparrow\downarrow\rangle + e^{i\varphi} \sin \vartheta |\downarrow\uparrow\rangle$$

Width of the range of J/k where the amplitude of AB oscillations is larger than 0.95 versus $\vartheta \in [0, 2\pi]$ and $\varphi \in [0, \pi]$



Physical explanation

- For $\phi = 0 \forall kL$ the two arms of the ring are perfectly symmetric $\Rightarrow \mathbf{S}_{12}^2$ is conserved
- $|\uparrow\rangle |\Psi^-\rangle \rightarrow |\uparrow\rangle |\Psi^-\rangle$: no spin-flip!
- When $\phi \neq 0$: $k_1 \neq k_2$ and symmetry is lost $\Rightarrow \mathbf{S}_{12}^2$ is not conserved

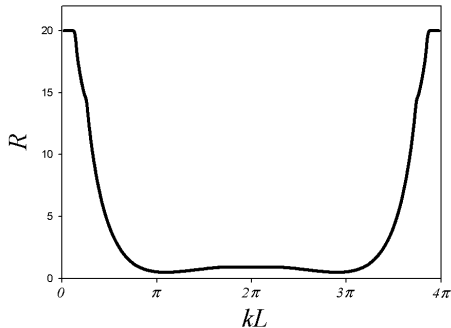
...however

- For $\phi \neq 0$ the s. m. of the wavefunction is periodic in kL with period 4π
- For $kL = 4n\pi$ the system behaves as if $kL = 0$ that is as if $k_1 = -k_2 = e\phi/\hbar cL \Rightarrow$ symmetry holds again $\Rightarrow \mathbf{S}_{12}^2$ is conserved



Quasi conservation of \mathbf{S}_{12}^2

Width R of the range of J/k where the amplitude of AB oscillations is larger than 0.95 vs. kL for the initial state $|\uparrow\rangle|\Psi^-\rangle$



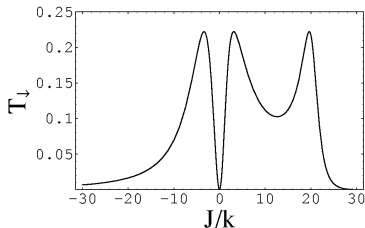
- R increases as $kL \rightarrow 0, 4\pi$
- Good robustness against deviations from $4n\pi$ ($R > 1$ for $kL < 0.8\pi$)
- Previous results shown in the case $kL = 1 \simeq 0.3\pi$



Scheme for generating $|\Psi^-\rangle$

At $\phi = 0$ strict conservation of \mathbf{S}_{12}^2 holds:

$$|\uparrow\rangle|\downarrow\downarrow\rangle \rightarrow t_\uparrow |\uparrow\rangle|\downarrow\downarrow\rangle + t_\downarrow |\downarrow\rangle|\downarrow\rangle |\Psi^+\rangle$$



- $kL = 1$
- One maximum for $J/k \simeq 3 \rightarrow$ within the nearly coherent range



Conclusions

- Magnetic impurities typically induce loss of electron coherence in mesoscopic devices
- Coherence can however be preserved when the impurities are allowed to be in suitable entangled states (cooperative effects)
- Quantum entanglement as an inhibitor of electron decoherence

