

DIVERGENCE OF THE ENTANGLEMENT RANGE IN LOW DIMENSIONAL QUANTUM SYSTEMS

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SUMMARY

- MODELS : ANTIFERROMAGNETIC EXCHANGE
vs UNIFORM MAGNETIC FIELD
- PHENOMENOLOGY : QUANTUM CRITICALITY, SATURATION,
FACTORIZATION
- TOOLS : CONCURRENCE
PARALLEL & ANTIPARALLEL ENTANGLEMENT
RANGE OF THE CONCURRENCE R
- RESULTS : ENTANGLEMENT "PHASE DIAGRAM"
DIVERGENCE OF R
- MORE MODELS
- CONCLUSIONS AND BIBLIOGRAPHY

MODELS

$$\mathcal{H} = J \sum_i S_i^x S_{i+1}^x + \Delta_y S_i^y S_{i+1}^y + \Delta_z S_i^z S_{i+1}^z - h S_i^z$$

$$J > 0 ; 0 \leq \Delta_{y,z} \leq 1 ; S = \frac{1}{2}$$

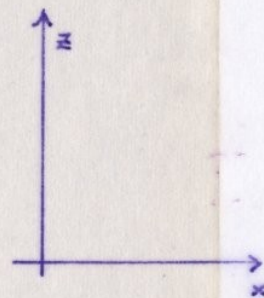
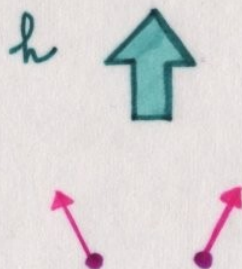
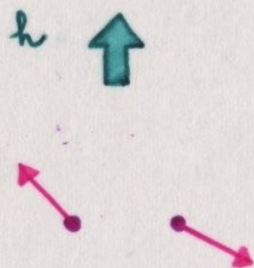
AN ANTIFERROMAGNETIC ZOOLOGY

Ising
XY(z)

Kosterlitz-Thouless
XX(z)

isotropic
XXX

WITH A COMMON FEATURE



ANTIFERROMAGNETIC EXCHANGE VERSUS UNIFORM MAGNETIC FIELD

// x

// z

PHENOMENOLOGY (T=0)

(TOOLS : $g_r^{\alpha\beta} \equiv \langle S_i^\alpha S_{i+r}^\beta \rangle$ & $M^\alpha \equiv \langle S_i^\alpha \rangle$; $\alpha, \beta = x, y, z$)

$M^z \neq 0$ whenever $h_z \neq 0$

• QUANTUM CRITICALITY [XY(Z) & XX(Z)]

$\exists h_c > 0$ such that

$$h < h_c \quad g_r^{xx} \sim \frac{1}{r^\eta} \quad \begin{cases} \eta = 0 \Rightarrow M^x > 0 \text{ ordered phase (Ising)} \\ \eta > 0 \Rightarrow M^x = 0 \text{ quasi-ordered phase (KT)} \end{cases}$$

$$h > h_c \quad g_r^{xx} \sim e^{-r/\xi^{xx}} \quad \text{disordered phase}$$

• SATURATION [XXX & XX(Z)]

$\exists h_s < +\infty$ such that

$$h \geq h_s \quad g_r^{zz} = (M^z)^2 = \frac{1}{4} \quad \forall r \quad \text{paramagnetic disordered phase}$$

FACTORIZATION

[XY(z) & XX(z)]

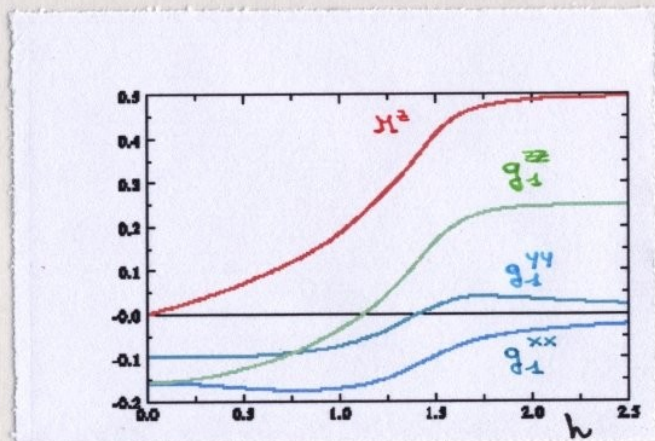
Kurmann, Thomas, and Müller : Physica 112A , 235 (1982)

$$\exists \omega_f = \sqrt{(1 + \Delta_z)(\Delta_y + \Delta_z)} \text{ such that}$$

$$|GS\rangle = \prod_i |\varphi_i\rangle \text{ "classical-like GS" } @ \omega_f$$

$$|\varphi_{2i}\rangle = |\alpha\rangle \text{ \& } |\varphi_{2i+1}\rangle = |\beta\rangle$$

$$\varepsilon_0 = -\frac{1 + \Delta_x + \Delta_y}{4}$$



no peculiarities @ ω_f

!

"new" TOOLS

ENTANGLEMENT BETWEEN A COUPLE OF SPINS

$$S_i \text{ \& \& } S_{i+r}$$

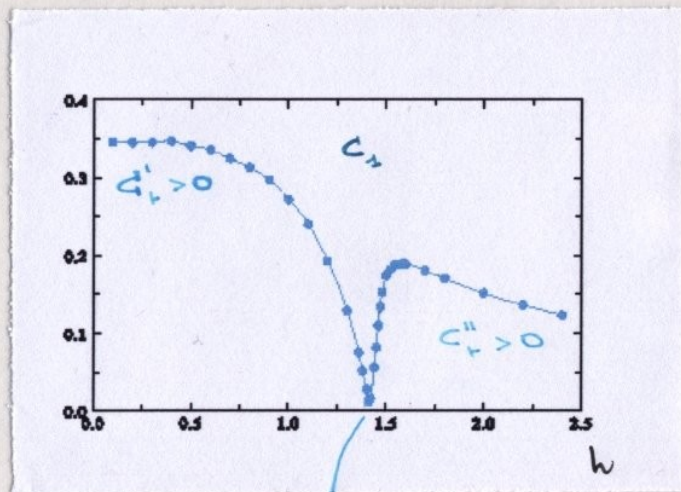


CONCURRENCE C_r

$$C_r \equiv 2 \max \{0, c'_r, c''_r\}$$

$$c'_r = |g_r^{xx} + g_r^{yy}| - \sqrt{\left(\frac{1}{4} + g_r^{zz}\right)^2 - H_z^2} ; \quad c''_r = |g_r^{xx} - g_r^{yy}| + g_r^{zz} - \frac{1}{4}$$

$$\begin{aligned} \underline{h < h_f} \\ c'_r &> 0 \\ c''_r &< 0 \end{aligned}$$



$$\begin{aligned} \underline{h > h_f} \\ c'_r &< 0 \\ c''_r &> 0 \end{aligned}$$

at $h_f \rightarrow c'_r = c''_r = 0$

PARALLEL & ANTIPARALLEL ENTANGLEMENT (bipartite)

$$|e_1\rangle \equiv \frac{|\uparrow\uparrow\rangle + |\downarrow\downarrow\rangle}{\sqrt{2}} ; |e_2\rangle \equiv \frac{|\uparrow\uparrow\rangle - |\downarrow\downarrow\rangle}{\sqrt{2}}$$

parallel Bell states

$$|e_3\rangle \equiv \frac{|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle}{\sqrt{2}} ; |e_4\rangle \equiv \frac{|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle}{\sqrt{2}}$$

antiparallel Bell states

$p_\alpha \Rightarrow$ probability that the pair S_i & S_{i+r} be in the state $|e_\alpha\rangle$

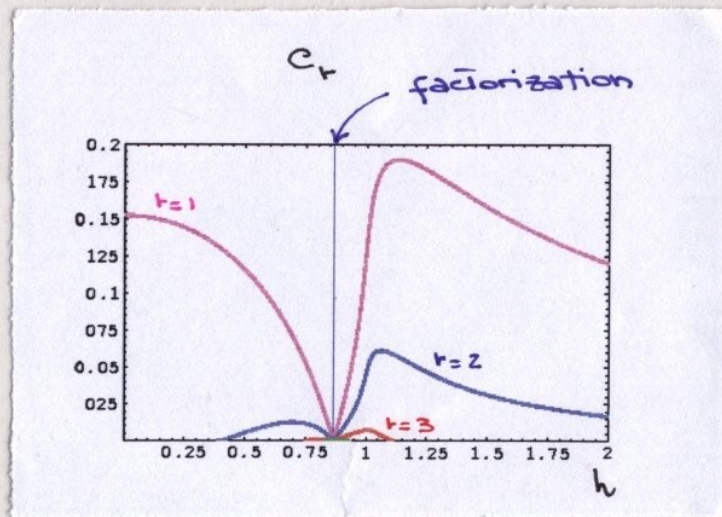
+ $p_{\uparrow\uparrow}$, $p_{\downarrow\downarrow}$

$$2C'_r = |p_3 - p_4| - 2\sqrt{p_{\uparrow\uparrow}p_{\downarrow\downarrow}}$$



$$2C''_r = 2\max\{p_1, p_2\} - 1$$

antiparallel
entanglement



parallel
entanglement

ENTANGLEMENT "PHASE DIAGRAM"

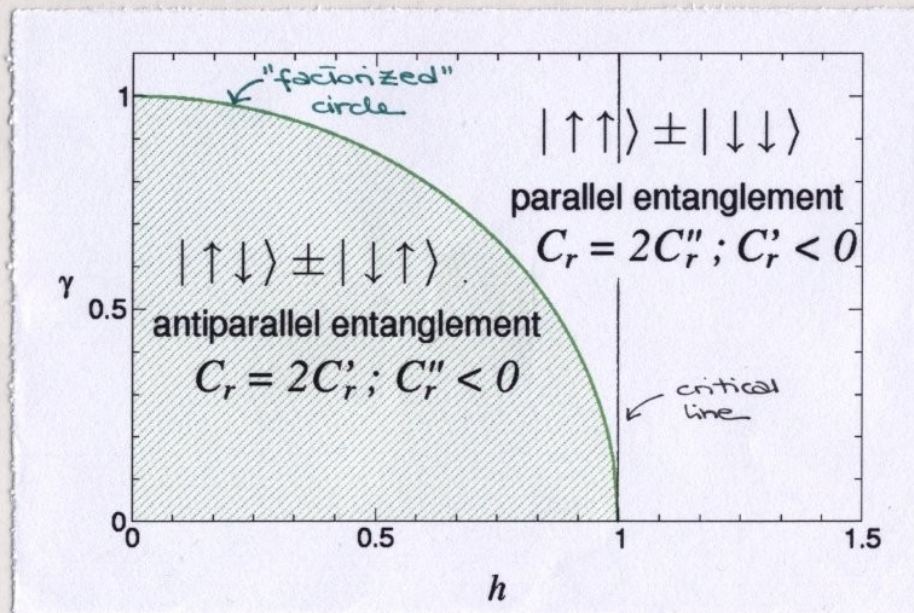
$$\left. \begin{array}{l} C_r' \geq 0 \\ C_r'' \leq 0 \end{array} \right\} h < h_f \quad \forall r !$$

$$C_r' = 0 = C_r'' \quad @ h_f \quad \forall r !$$

$$\left. \begin{array}{l} C_r' \leq 0 \\ C_r'' \geq 0 \end{array} \right\} h > h_f \quad \forall r !$$



XY model
in a transverse
field

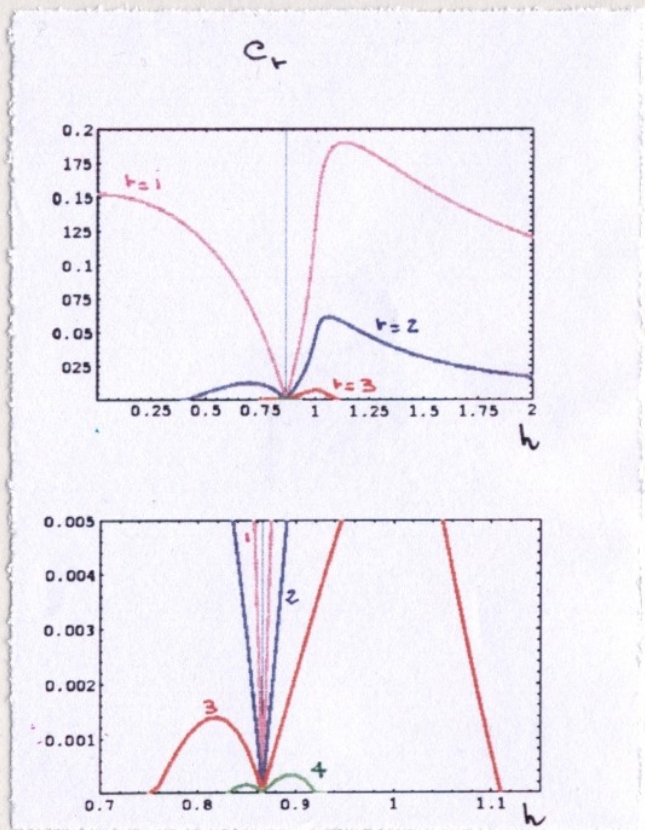
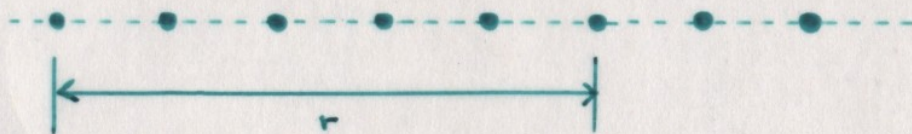


$$h_f = \sqrt{1-\gamma^2}$$

$$h_c = 1$$

?

along the chain



as $h \rightarrow h_f$
all C_r get positive



long-ranged bipartite
entanglement

but

$$\sum_r C_r^2 \leq \tau \leq 1$$



$$C_r \rightarrow 0$$

$$h \rightarrow h_f$$

$$\forall r$$

ENTANGLEMENT RANGE

it is observed that

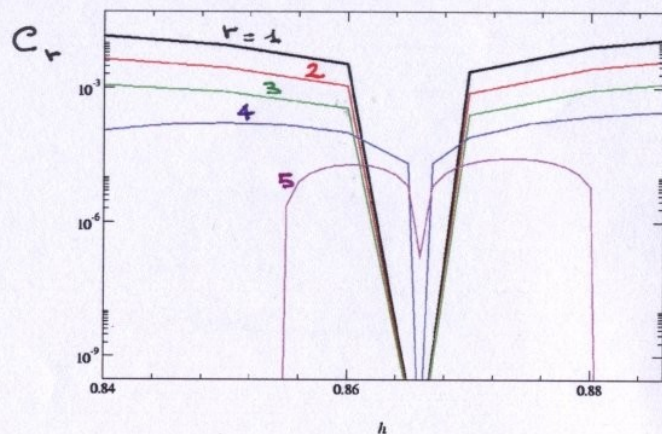
$$\text{if } C_r > 0 \Rightarrow C_m > 0 \quad \forall m \leq r$$

the range of the concurrence may be hence defined as

$$R \text{ such that } \begin{array}{ll} C_r > 0 & \text{if } r \leq R \\ C_r < 0 & \text{if } r > R \end{array}$$

for a given model R depends on h

DIVERGENCE of R @ h_f



$$\lim_{h \rightarrow h_f} R = \infty$$

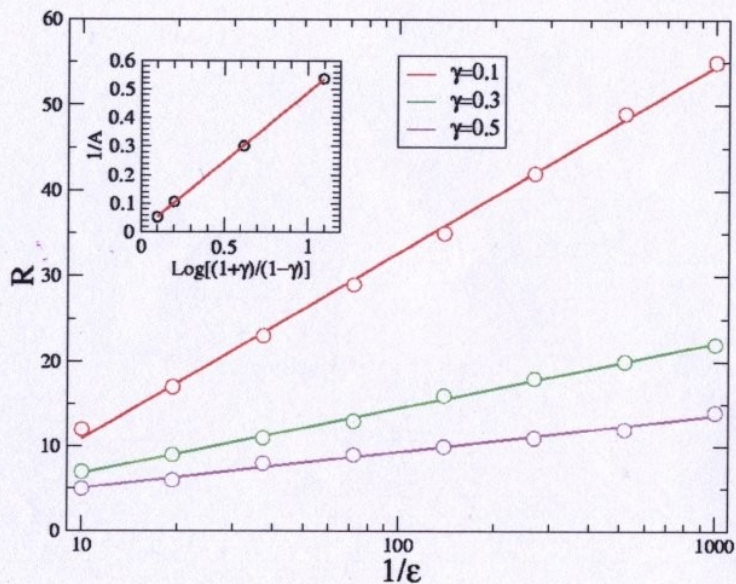
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the way R diverges @ h_f
depends on the model:

$$R^{xy} \sim \frac{1}{\ln\left(\frac{1-\gamma}{1+\gamma}\right)} \frac{1}{\varepsilon} = \frac{1}{\ln(\Delta\gamma)} \frac{1}{\varepsilon}$$

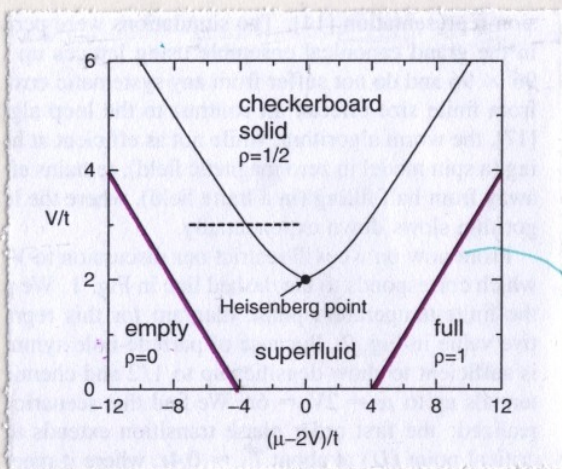
$$\varepsilon \equiv |h - h_f| \quad ; \quad h_f = \sqrt{1 - \gamma^2} = \sqrt{\Delta\gamma(1 + \Delta\gamma)}$$

$$(A' \equiv \varepsilon R^{xy} \sim \ln(\Delta\gamma))$$



MORE MODELS

- factorization occurs in any dimension provided the lattice is bipartite
→ (ladders, square lattices ...)
- factorization is observed in magnetic models with long-ranged interactions
(Vidal et al. PRL 90, 227902 (2003))
- XXZ models in a field are exactly mapped into the
hard-core boson Hubbard models



H-c boson Hubbard model
on the square lattice

$$\mathcal{H} = -t \sum_{\langle ij \rangle} (a_i^\dagger a_j + a_j^\dagger a_i) + V \sum_{\langle ij \rangle} n_i n_j - \mu \sum_i n_i$$

→ XXZ via $t \rightarrow \frac{1}{2}$; $V \rightarrow iz$; $\mu - 2V \rightarrow h$

— $\sim h = 2(1 + iz)$ saturation in 2d

...

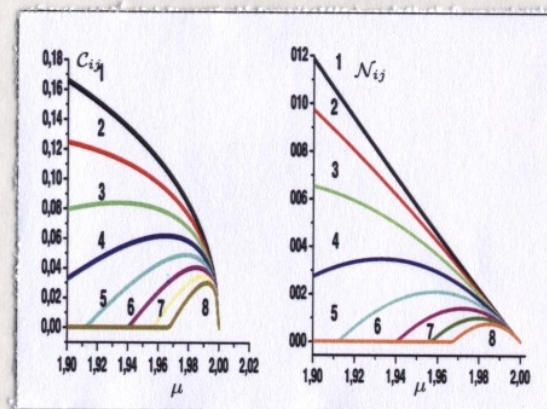
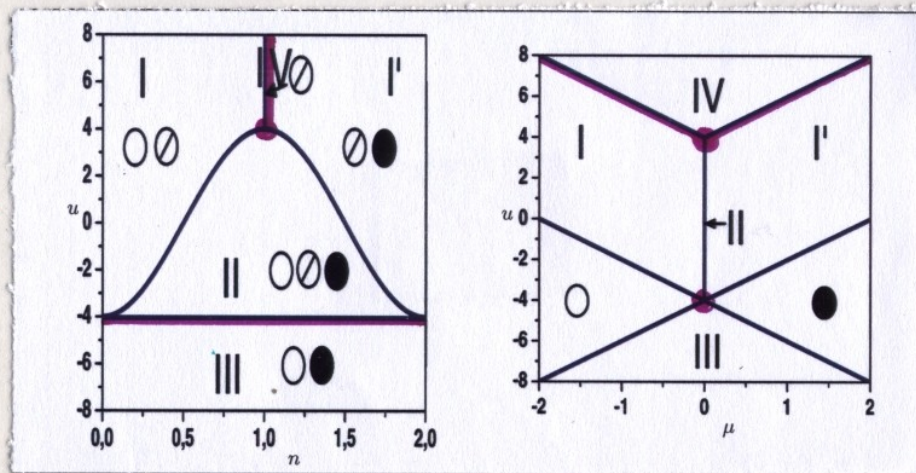
$T > 0$?

EVEN MORE MODELS

the bond-charge extended Hubbard model

$$\mathcal{H}_{BC} = - \sum_{i,\epsilon} [1 - \kappa (n_{i,\bar{\epsilon}} + n_{i+1,\bar{\epsilon}})] c_{i,\epsilon}^{\dagger} c_{i+1,\epsilon} - \mu \sum_{i,\epsilon} n_{i,\epsilon} + u \sum_i (n_{i,\uparrow} - \frac{1}{2})(n_{i,\downarrow} - \frac{1}{2})$$

$c_{i,\epsilon}^{\dagger}, c_{i,\epsilon}$: fermionic operators ; $\epsilon = \uparrow, \downarrow$; $\bar{\epsilon}$ opposite of ϵ ; $n_{i,\epsilon} = c_{i,\epsilon}^{\dagger} c_{i,\epsilon}$



$I' \rightarrow IV$

κ — long-ranged bipartite entanglement $\longrightarrow$ divergent R



FACTORIZATION $\approx$ QUANTUM CRITICALITY



studied models (always with a uniform field $\parallel \hat{z}$)

1d : XY

XYX

XXX

with $\Delta_y = 0, 0.25, 0.5, 0.75$

ladder : XXX

2d : XYX

with $\Delta_y = 0, 0.25, 0.5, 0.75$

Phys. Rev. Lett. 93, 167203 (2004)

Phys. Rev. Lett. 94, 147208 (2005)

Eur. Phys. J. D38, 563 (2006)

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THANKS !